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The Crystal Structure of Methyltriphenylphosphonium Bis-7,7,8,8-tetracyanoquinodimethanide at 53°C

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The crystal structure of the high-temperature modification of $(\text{CH}_3\text{PPh}_3)^+(\text{TCNQ})_2^-$ has been determined at 53°C. The crystals are triclinic with the space group $P\bar{1}$ and lattice constants: $a = 15.780$ (5), $b = 14.545$ (4), $c = 8.910$ (3) Å, $\alpha = 117.22$ (4), $\beta = 98.05$ (7), $\gamma = 97.02$ (6)° and $Z = 2$. The structure was refined by the block-diagonal least-squares method with 2514 reflexions measured on a four-circle manual diffractometer to give an R value of 0.051. The structure is closely related to that of the low-temperature modification, consisting of tetrads of TCNQ molecules and methyltriphenylphosphonium ions. The tetrads are arranged to form a columnar structure along the b axis. The average interplanar distances in the tetrads increased to 3.32 and 3.26 Å, which are longer by 0.08 and 0.04 Å, respectively, than the corresponding distances in the low-temperature modification, whereas the interplanar distance between the two adjacent tetrads of 3.55 Å was shortened by about 0.02 Å. The overlapping mode of the TCNQ moieties within a tetrad is unaltered. The nearest neighbour overlap between the two adjacent tetrads changed from a modified ring-ring overlap to an overlap involving a diagonal shift of the molecular centres. The conformation of $(\text{CH}_3\text{PPh}_3)^+$ ion is different from that of the low-temperature form; two of the phenyl groups are rotated by about 56° and 45° from the positions observed in the low-temperature form. Moreover, the conformation of the phosphonium ion is disordered, two of the phenyl rings taking different orientations with a probability of about 0.1. These changes of the structure are in accord with the change in physical properties of the high-temperature phase.

Introduction

Crystals of $(\text{CH}_3\text{PPh}_3)^+(\text{TCNQ})_2^-$ undergo a first-order phase transition at 315.7 K (Kepler, 1963). The temperature dependence of physical properties such as electrical conductivity, magnetic susceptibility etc., has been extensively studied by a number of workers (for example, Jones & Chesnut, 1963; Iida, Kinoshita, Sano & Akamatu, 1964; Iida, Kinoshita, Kawamori & Suzuki, 1964; Nordio, Soos & McConnell, 1966; Kosaki, Iida, Sorai, Suga & Seki, 1970; Shirotani, Kawamura & Iida, 1972). Models for this transition were also proposed (Chesnut, 1964; Merkl, Hughes, Berliner & McConnell, 1965). The crystal structure of this compound at room temperature has been determined by the photographic method (McPhail, Seme-

niuk & Chesnut, 1971). Quite independently we have determined the crystal structure by the counter method (Tsuchiya, Marumo & Saito, 1973). Positions of hydrogen atoms were also determined. The final R value was 0.060 for 3574 observed reflexions. Both results agree well, although the cell dimensions are somewhat different. The crystal structure of the high-temperature modification was determined in order to gain a better understanding of the phase transition. The result will be discussed with particular reference to the structure of the low-temperature phase.

Experimental

The crystals of $(\text{CH}_3\text{PPh}_3)^+(\text{TCNQ})_2^-$ were kindly supplied by Dr Iida of Hokkaido University. A single-

THE CRYSTAL STRUCTURE OF $(CH_3PPh_3)^+(TCNQ)_2$ AT 53°C

crystal specimen was found to transform into the high-temperature form without breaking, if the temperature was raised slowly at a rate of 0.4° per min. A part of the

reciprocal-lattice network containing those reflexions of which the intensities change sensitively with the phase transition was recorded on Weissenberg photo-

Table 1. Observed and calculated structure factors

M	FO	FC	M	FO	FC	M	FO	FC	M	FO	FC	M	FO	FC	M	FO	FC	M	FO	FC	M	FO	FC	M	FO	FC	M	FO	FC																																																																																																																																																																																																																																																																																									
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Table 1 (*cont.*)

graphs at different temperatures. It was confirmed that the phase change was completed in a very short time by this heating rate and no indication was observed of the low-temperature phase migrating into the high-temperature phase. The crystal specimen was heated by a stream of electrically heated air. Thermocouple measurements indicated that the temperature was constant at $53 \pm 1^\circ\text{C}$ throughout the experiment. Intensity data were collected on a manual four-circle diffractometer. The crystals disintegrated during exposure to X-rays and measurements had to be made with half a dozen specimens, whose maximum dimensions ranged from 0.3 to 0.4 mm and which were cut from larger crystals. Mo $K\alpha$ radiation was used with the ω -2 θ scan technique. Data were recorded for 3250 independent reflexions up to $2\theta=40^\circ$, of which 2514 with $|F| \geq 3\sigma$, were regarded as 'observed'. A standard reflexion was measured every two hours. The intensities were corrected for Lorentz and polarization effects, but corrections for absorption and extinction were not applied. The crystal data are: $(\text{C}_{19}\text{H}_{18}\text{P}) (\text{C}_{12}\text{H}_4\text{N}_4)_2$, M.W. 685.7. Triclinic, $a = 15.780$ (5), $b = 14.545$ (4), $c = 8.910(3)$ Å, $\alpha = 117.22$ (4), $\beta = 98.05$ (7), $\gamma = 97.02$ (6)°; $U = 1759.9$ Å 3 , $Z = 2$, $D_x = 1.299$ g cm $^{-3}$, at 53°C , Mo $K\alpha$ ($\lambda = 0.7107$ Å), $\mu = 1.28$ cm $^{-1}$. Space group $\overline{P}1$.

Crystal data of the low-temperature modification are: Triclinic, $a = 15.762$ (3), $b = 14.725$ (5), $c = 9.022$ (4) Å, $\alpha = 120.01$ (3), $\beta = 97.32$ (6), $\gamma = 96.21$ (7)°, $U = 1762.1$ Å³. (Tsuchiya, Marumo & Saito, 1973). Space group $P\bar{1}$. Triclinic, $a = 15.74$, $b = 14.67$, $c = 9.01$ Å, $\alpha = 119.98$, $\beta = 97.04$, $\gamma = 96.55$ °, $U = 1751$ Å³ (McPhail *et al.*, 1971).

McPhail's data (a' , b' , c') were transformed according to the relations: $a = b' + c'$, $b = -a' - b'$, $c = a'$.

Structure determination

Initial positional parameters for the non-hydrogen atoms were taken from the final parameters of the low-temperature modification (Tsuchiya *et al.*, 1973), since the general features of the intensity distribution indicated that the structure is closely related to that of the low-temperature modification. The conventional *R* value was 0.35 at this stage. The atomic positions and thermal parameters were refined in the usual way by the least-squares method with a block-diagonal program *HBLS-4* written by Dr Ashida. With the introduction of anisotropic temperature factors further cycles of least-squares refinement reduced the *R* index to 0.13. A difference Fourier synthesis showed all the hydrogen atoms in plausible positions. Further refinement cycles were carried out with anisotropic temperature factors for non-hydrogen atoms and with isotropic temperature factors for hydrogen atoms. The *R* value converged to 0.070. In a difference Fourier synthesis at this stage, however, there were several well defined peaks with heights of $0.4 \text{ e } \text{\AA}^{-3}$. The maps can be explained if two of the phenyl rings are considered to assume different orientations with a population of about one tenth. This disorder was taken into account and the populations of the carbon atoms were also refined. Isotropic temperature factors were assumed for these carbon atoms. Positions of hydrogen atoms attached to the carbon

atoms in this less populated position could not be determined. After several refinement cycles the structure converged with $R=0.051$ and $R_2=0.048$. At the final stage of the refinement all the parameter shifts were less than one eighth of the corresponding standard deviations. Unit weight was given to all the reflexions. The atomic scattering factors were taken from *International Tables for X-ray Crystallography* (1962). Table 1 gives the observed and calculated structure factors. The atomic parameters and their standard deviations are listed in Table 2.

Thermal motion

The motion of two TCNQ molecules has been analysed in terms of the rigid-body tensors of translation (\mathbf{T}) and libration (ω) (Cruickshank, 1965). Rigid-body parameters for the two quinodimethane moieties are listed in Table 3. There are two crystallographically independent TCNQ moieties, hereafter designated as TCNQ I and II, respectively. In both molecules, the translation is quite isotropic, whereas the libration is quite anisotropic and the axis of the largest r.m.s. amplitude corresponds to the long axis of the molecule

Table 2. *Atomic parameters of non-hydrogen atoms ($\times 10^4$)*

The values of β_{ij} refer to the expression: $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$.

Here and elsewhere in this paper the estimated standard deviations in the last figure are given in parentheses.

	x	y	z	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}	Population
P	4992 (1)	2989 (1)	1541 (2)	34 (1)	58 (1)	120 (2)	13 (1)	13 (1)	39 (1)	
C(1)	4303 (4)	3495 (4)	3048 (7)	53 (3)	89 (5)	163 (11)	20 (3)	35 (5)	24 (6)	
C(2)	4376 (3)	2434 (4)	-607 (6)	30 (2)	57 (4)	174 (10)	8 (2)	-7 (4)	54 (5)	
C(3)	4800 (4)	2092 (5)	-1973 (6)	66 (4)	122 (6)	115 (11)	13 (4)	29 (5)	62 (7)	
C(4)	4289 (5)	1622 (6)	-3645 (7)	126 (6)	137 (7)	127 (12)	21 (5)	37 (7)	65 (8)	
C(5)	3395 (4)	1534 (5)	-3904 (8)	87 (4)	118 (6)	265 (15)	-7 (4)	-82 (7)	101 (8)	
C(6)	2992 (4)	1829 (5)	-2585 (8)	65 (4)	92 (6)	320 (16)	11 (4)	-12 (6)	98 (8)	
C(7)	3475 (3)	2288 (4)	-937 (7)	50 (3)	77 (5)	182 (11)	13 (3)	-14 (5)	50 (6)	
C(8)	5445 (3)	2010 (4)	1828 (6)	28 (3)	56 (4)	103 (10)	5 (3)	13 (4)	39 (5)	0.9
C(9)	4927 (4)	1339 (5)	2189 (8)	65 (4)	84 (6)	265 (15)	34 (4)	70 (6)	93 (8)	0.9
C(10)	5275 (5)	559 (5)	2409 (9)	95 (5)	88 (6)	272 (16)	28 (4)	58 (7)	108 (9)	0.9
C(11)	6136 (4)	485 (5)	2315 (8)	61 (4)	97 (6)	257 (15)	29 (4)	27 (6)	101 (8)	0.9
C(12)	6614 (4)	1110 (5)	1869 (10)	42 (3)	110 (6)	390 (20)	28 (4)	46 (7)	139 (10)	0.9
C(13)	6275 (4)	1884 (5)	1649 (9)	48 (4)	89 (6)	327 (17)	23 (4)	54 (6)	121 (9)	0.9
C(14)	5855 (3)	4063 (4)	1986 (7)	33 (3)	56 (4)	161 (11)	10 (3)	1 (5)	50 (6)	0.9
C(15)	6487 (4)	4510 (5)	3450 (8)	44 (3)	101 (6)	204 (14)	-4 (4)	-17 (5)	102 (8)	0.9
C(16)	7146 (4)	5346 (5)	3841 (8)	53 (4)	109 (7)	236 (15)	-9 (4)	-26 (6)	87 (9)	0.9
C(17)	7201 (4)	5759 (5)	2785 (9)	50 (4)	80 (6)	306 (17)	-12 (4)	-14 (6)	70 (9)	0.9
C(18)	6597 (5)	5288 (5)	1207 (9)	82 (5)	72 (6)	302 (17)	1 (4)	17 (7)	99 (9)	0.9
C(19)	5915 (4)	4457 (5)	801 (8)	56 (4)	74 (5)	204 (14)	2 (4)	-6 (6)	62 (7)	0.9
C(20)	-955 (3)	1131 (4)	-1640 (5)	32 (2)	49 (4)	101 (9)	11 (2)	11 (4)	33 (5)	
C(21)	-1100 (3)	1331 (4)	31 (6)	33 (2)	55 (4)	123 (9)	12 (2)	22 (4)	40 (5)	
C(22)	-98 (3)	1020 (4)	-1939 (5)	38 (3)	64 (4)	93 (9)	13 (3)	19 (4)	37 (5)	
C(23)	-1631 (3)	1055 (4)	-2905 (6)	37 (3)	60 (4)	117 (9)	10 (3)	22 (4)	40 (5)	
C(24)	-2483 (3)	1168 (4)	-2632 (6)	40 (3)	81 (5)	144 (10)	7 (3)	7 (4)	48 (6)	
C(25)	-1505 (3)	860 (4)	-4572 (6)	37 (3)	69 (4)	147 (10)	9 (3)	4 (4)	45 (6)	
N(26)	-3168 (3)	1256 (4)	-2437 (6)	43 (3)	153 (6)	280 (12)	26 (3)	28 (5)	106 (7)	
N(27)	-1433 (3)	690 (4)	-5927 (6)	66 (3)	129 (5)	168 (10)	24 (3)	24 (4)	79 (6)	
C(28)	410 (3)	1312 (3)	975 (6)	34 (2)	44 (4)	111 (9)	10 (2)	10 (4)	33 (5)	
C(29)	555 (3)	1106 (4)	-708 (6)	34 (3)	56 (4)	114 (9)	10 (3)	26 (4)	36 (5)	
C(30)	-448 (3)	1402 (4)	1248 (6)	40 (3)	62 (4)	102 (9)	12 (3)	18 (4)	43 (5)	
C(31)	1086 (3)	1408 (4)	2255 (6)	32 (2)	61 (4)	120 (9)	8 (3)	16 (4)	44 (5)	
C(32)	1940 (3)	1327 (4)	1981 (6)	41 (3)	70 (4)	154 (10)	7 (3)	1 (4)	59 (6)	
C(33)	976 (3)	1610 (4)	3933 (6)	32 (3)	68 (4)	155 (10)	9 (3)	-4 (4)	51 (6)	
N(34)	2642 (3)	1277 (4)	1796 (6)	37 (2)	133 (5)	299 (12)	17 (3)	16 (4)	117 (7)	
N(35)	913 (3)	1782 (4)	5291 (5)	61 (3)	123 (5)	159 (9)	20 (3)	14 (4)	79 (6)	
C(36)	-1386 (3)	3564 (4)	-2050 (6)	37 (3)	58 (4)	128 (10)	16 (3)	21 (4)	46 (5)	
C(37)	-1547 (3)	3739 (4)	-391 (6)	31 (3)	67 (4)	142 (10)	17 (3)	28 (4)	42 (5)	
C(38)	-520 (3)	3487 (4)	-2304 (6)	42 (3)	74 (4)	143 (10)	17 (3)	28 (4)	67 (6)	
C(39)	-2036 (3)	3504 (4)	-3325 (6)	40 (3)	77 (4)	162 (11)	16 (3)	23 (4)	69 (6)	
C(40)	-2914 (4)	3591 (5)	-3169 (6)	55 (3)	110 (5)	151 (11)	28 (3)	28 (5)	85 (7)	
C(41)	-1865 (3)	3364 (4)	-4936 (6)	40 (3)	99 (5)	166 (11)	14 (3)	7 (4)	81 (6)	
N(42)	-3618 (3)	3682 (5)	-3087 (7)	60 (3)	202 (7)	313 (13)	49 (4)	55 (5)	156 (8)	
N(43)	-1734 (3)	3262 (4)	-6219 (6)	54 (3)	172 (6)	195 (10)	15 (3)	20 (4)	118 (7)	
C(44)	-49 (3)	3728 (4)	580 (6)	40 (3)	49 (4)	115 (9)	14 (2)	22 (4)	42 (5)	
C(45)	113 (3)	3567 (4)	-1068 (6)	32 (3)	65 (4)	166 (10)	12 (3)	19 (4)	61 (6)	
C(46)	-915 (3)	3808 (4)	844 (6)	40 (3)	66 (4)	137 (10)	17 (3)	26 (4)	57 (5)	
C(47)	621 (3)	3799 (4)	1832 (6)	33 (3)	47 (4)	146 (10)	12 (2)	18 (4)	35 (5)	
C(48)	1492 (3)	3762 (4)	1583 (6)	41 (3)	71 (4)	164 (11)	13 (3)	8 (4)	49 (6)	
C(49)	505 (3)	3918 (4)	3475 (6)	35 (3)	56 (4)	136 (10)	12 (3)	3 (4)	33 (5)	
N(50)	2190 (3)	3741 (4)	1423 (7)	43 (3)	131 (5)	311 (13)	27 (3)	34 (5)	97 (7)	
N(51)	459 (3)	4005 (4)	4795 (5)	61 (3)	94 (4)	151 (9)	19 (3)	21 (4)	55 (5)	

Table 2 (cont.)

				B^* (\AA^2)	Popula-tion
C'(8)	5534 (31)	1957 (37)	1786 (60)	4·4 (1·1)	0·1
C'(9)	5566 (22)	1895 (27)	3317 (43)	1·2 (0·7)	0·1
C'(10)	6008 (26)	1226 (32)	3464 (51)	2·4 (0·8)	0·1
C'(11)	6294 (33)	491 (40)	2040 (65)	5·3 (1·1)	0·1
C'(12)	6247 (32)	563 (39)	325 (63)	4·6 (1·1)	0·1
C'(13)	5782 (24)	1329 (30)	359 (48)	2·1 (0·8)	0·1
C'(14)	5775 (33)	4076 (40)	2003 (64)	5·0 (1·1)	0·1
C'(15)	6635 (38)	3982 (46)	2245 (75)	6·7 (1·4)	0·1
C'(16)	7179 (36)	4912 (44)	2670 (71)	6·4 (1·3)	0·1
C'(17)	6912 (37)	5752 (44)	2794 (73)	6·1 (1·3)	0·1
C'(18)	6063 (28)	5861 (35)	2679 (56)	3·7 (0·9)	0·1
C'(19)	5490 (20)	4912 (25)	2174 (39)	0·9 (0·6)	0·1

* Isotropic temperature factor

Positional parameters for the hydrogen atoms ($\times 10^3$)

	x	y	z	Population
H(C3)	548 (3)	218 (4)	-176 (6)	
H(C4)	453 (4)	141 (5)	-457 (8)	
H(C5)	301 (3)	123 (4)	-498 (7)	
H(C6)	234 (3)	173 (4)	-274 (6)	
H(C7)	322 (3)	255 (3)	26 (5)	
H(C9)	431 (4)	148 (5)	239 (8)	0·9
H(C10)	488 (5)	-2 (6)	265 (9)	0·9
H(C11)	637 (4)	-18 (5)	238 (7)	0·9
H(C12)	725 (4)	103 (5)	179 (8)	0·9
H(C13)	665 (4)	237 (5)	136 (8)	0·9
H(C15)	645 (3)	433 (4)	435 (6)	0·9
H(C16)	752 (4)	557 (4)	490 (7)	0·9
H(C17)	766 (4)	636 (5)	311 (8)	0·9
H(C18)	662 (4)	556 (4)	25 (8)	0·9
H(C19)	548 (4)	409 (4)	-39 (7)	0·9
Ha(Cl)	457 (4)	384 (5)	428 (8)	
Hb(Cl)	382 (4)	285 (4)	277 (7)	
Hc(Cl)	398 (5)	406 (6)	277 (9)	
H(C21)	-173 (3)	141 (4)	21 (6)	
H(C22)	-4 (3)	91 (4)	-307 (6)	
H(C29)	117 (4)	98 (4)	-93 (7)	
H(C30)	-53 (4)	147 (5)	241 (7)	
H(C37)	-214 (4)	380 (5)	-33 (8)	
H(C38)	-45 (4)	327 (5)	-361 (7)	
H(C45)	71 (4)	345 (4)	-118 (7)	
H(C46)	-104 (4)	385 (5)	194 (8)	

Mean isotropic temperature factor of hydrogen atom is

 $3\cdot5$ (1·4) \AA^2

for which the moment of inertia is the smallest. A similar analysis was carried out for the entire molecules; however, the r.m.s. deviation ΔU_{ij} between observed and calculated thermal parameters increased by a factor of two. This poor agreement is probably due to the wagging of the cyano groups. *ORTEP* plots of the thermal-motion ellipsoids for the $(\text{CH}_3\text{PPh}_3)^+$ ion of the low and high-temperature phases are illustrated in Fig. 1 (Johnson, 1965). A rigid-body analysis for an entire $(\text{CH}_3\text{PPh}_3)^+$ ion gave a r.m.s. deviation ΔU_{ij} of $0\cdot0128 \text{\AA}^2$; thus the validity of using the rigid-body approximation is questionable. As discerned from Fig. 1, each phenyl ring exhibits the largest librational motion around the P-C bond. The amplitudes of the thermal motions increase in the high-temperature phase whereas those of the phenyl ring II in the low-temperature phase were larger compared with those of phenyl ring I and III and these large amplitudes did not alter appreciably in the high-temperature phase.

Description of the structure and discussion

The structure viewed along the b axis is illustrated in Fig. 2. It consists of columns of TCNQ's packed almost parallel to each other. Within the columns, which run parallel to the b axis, a tetrad of molecules may be recognized. The columns are packed side by side to form layers parallel to (100). The $(\text{CH}_3\text{PPh}_3)^+$ ions are arranged between the layers of TCNQ columns. The average interplanar spacings between TCNQ I...II and between I...I' in a tetrad are 3·26 and 3·32 \AA respectively, where the prime indicates the molecule related to the first by a centre of symmetry. Between the tetrads the interplanar spacing II...II' is 3·55 \AA . Fig. 3 shows the dimensions of TCNQ I and II. Deviations of the atoms from the least-squares plane through the quinodimethane moiety are also indicated. The standard deviations of atomic coordinates are not small enough to determine from these dimensions whether the TCNQ's occur as TCNQ^0 and TCNQ^- or whether they are really indistinguishable as $\text{TCNQ}^{1/2-}$. A

Table 3. Rigid-body thermal parameters of TCNQ (quinodimethane moiety)

(a) Principal axes of the molecules in the form: $L\alpha + M\beta + N\gamma$

Moment of inertia (A.W. $\times \text{\AA}^2$)	L	TCNQ I	N
71·9	0·0484	0·0063	0·0918
259·2	-0·0409	0·0124	0·0803
331·1	0·0167	0·0779	0·0433

Moment of inertia (A.W. $\times \text{\AA}^2$)	L	TCNQ II	M	N
72·2	0·0474	0·0054	0·0926	
256·7	-0·0422	0·0100	0·0774	
328·9	0·0161	0·0783	0·0469	

(b) Molecular vibrational tensors

$T \times 10^4$	TCNQ I			TCNQ II		
	382	41	8	$T \times 10^4$	415	-10
		329	24\AA^2		354	49
			345			0\AA^2
$\omega \times 10$	298	-66	-10		284	412
		62	14 ($^\circ$) ²		-32	20
			15		37	6 ($^\circ$) ²
						32

Table 3 (cont.)

(c) Principal axes of the \mathbf{T} and $\boldsymbol{\omega}$ tensors relative to the molecular axes.

R.m.s. amplitude	TCNQ I			R.m.s. amplitude	TCNQ II		
	Direction cosines				Direction cosines		
0.172 Å	0.409	0.799	0.442	0.187 Å	0.350	0.895	-0.276
0.189	0.285	0.348	-0.893	0.192	0.599	-0.440	-0.669
0.201	0.867	-0.491	0.085	0.215	0.721	-0.069	0.690
0.88°	0.116	0.373	0.920	1.51°	-0.134	-0.627	0.767
2.26	-0.223	-0.893	0.390	2.02	-0.048	-0.769	-0.637
5.62	0.968	-0.251	-0.020	5.38	0.990	-0.123	0.073

(d) R.m.s. deviation between obs. and calc. U_{ij}

$$26 \times 10^{-4} \text{ Å}^2$$

$$38 \times 10^{-4} \text{ Å}^2$$

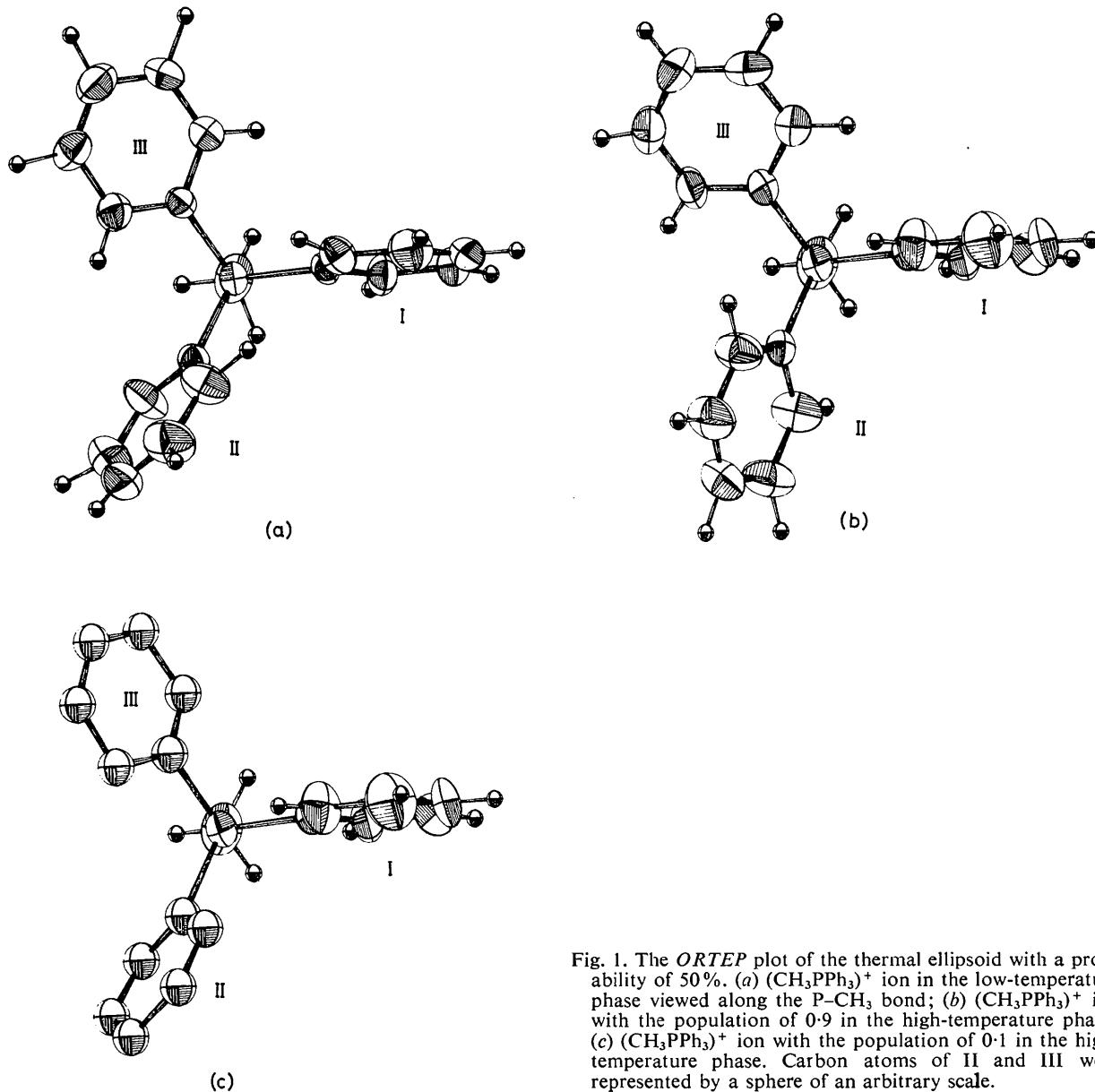


Fig. 1. The *ORTEP* plot of the thermal ellipsoid with a probability of 50%. (a) $(\text{CH}_3\text{PPh}_3)^+$ ion in the low-temperature phase viewed along the P-CH₃ bond; (b) $(\text{CH}_3\text{PPh}_3)^+$ ion with the population of 0.9 in the high-temperature phase; (c) $(\text{CH}_3\text{PPh}_3)^+$ ion with the population of 0.1 in the high-temperature phase. Carbon atoms of II and III were represented by a sphere of an arbitrary scale.

multivariate significance test (Cruickshank & Robertson, 1953) of the differences in dimensions between the two TCNQ moieties gave a value of $T^2=7$ which is not significant for four degrees of freedom. Within a tetrad there is one type of overlap between adjacent TCNQ units: ring-external bond type. This overlapping mode is the same as that observed in the low-temperature modification. However, the nearest neighbour overlap between adjacent tetrads changed slightly as shown in Fig. 4. Above the transition temperature it is an overlap involving diagonal shift of molecular centres similar to that observed in *N*-(*n*-propyl)quino-linium TCNQ₂ (Sundaresan & Wallwork, 1972). The least-squares planes through the TCNQ's indicate that they are not quite planar but bowed slightly; in molecule I in particular the terminal cyano groups are twisted slightly around the C=C bond. These features are the same as those observed in the low-temperature modification (McPhail *et al.*, 1971).

Bond distances and angles for the cation are listed in Table 4 with their estimated standard deviations.

As discerned from Fig. 1, the configuration of the (CH_3PPh_3)⁺ ion is quite different from a symmetrical propeller shape in both modifications. In the low-temperature phase, one phenyl ring (I) is nearly parallel to the P-CH₃ bond and the other two rings, II and III, are inclined by about -19 and -67°, respectively with respect to the P-CH₃ bond. In the high-temperature

phase, the cation takes two different configurations with different probability. In one configuration with larger population (0.9), phenyl ring I is almost parallel to the P-CH₃ bond as in the case of the low-temperature phase; however, rings II and III are inclined at

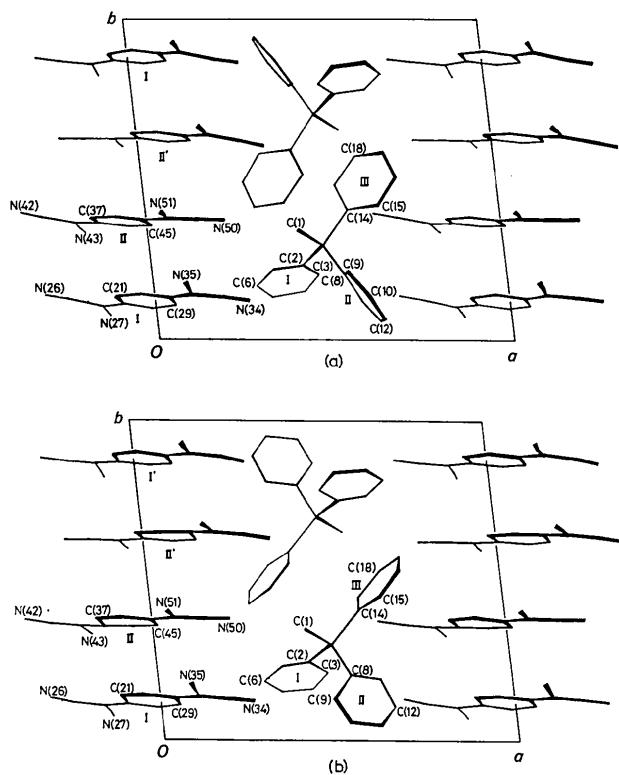


Fig. 2. The structure viewed along the b axis.
(a) Low-temperature phase. (b) High-temperature phase.

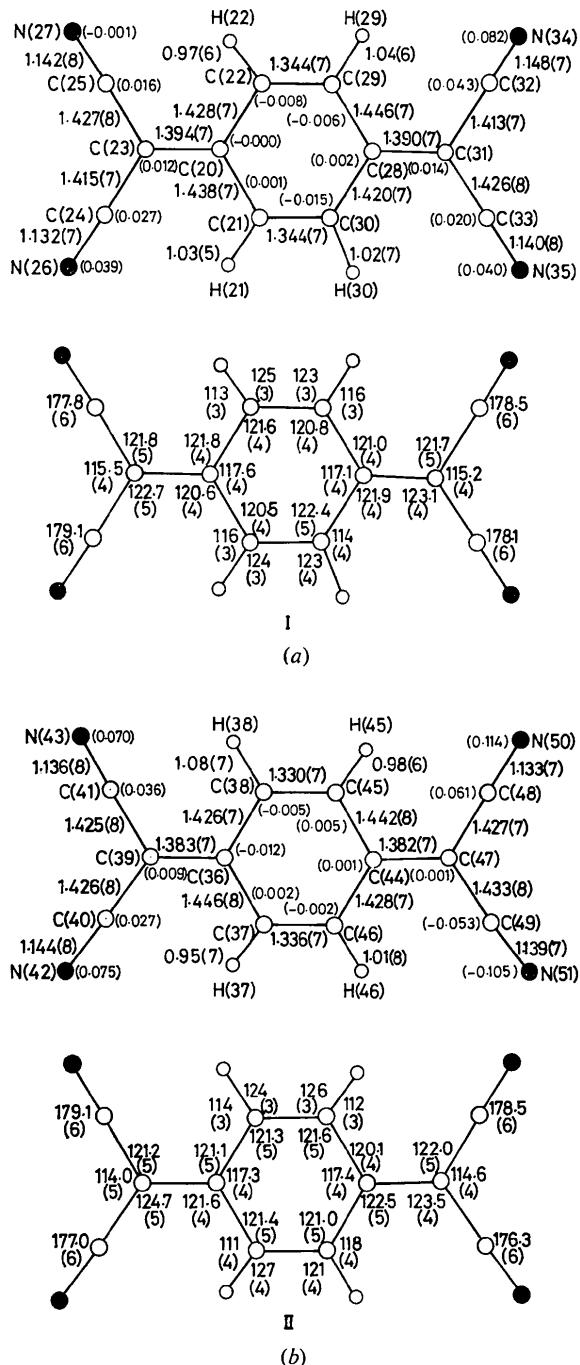


Table 4. Bond distances and angles for $(CH_3PPh_3)^+$

P····C(1)	1.797 (6) Å	C(18)···C(19)	1.387 (9) Å
P····C(2)	1.770 (5)	C(19)···C(14)	1.421 (11)
P····C(8)	1.769 (6)	C(1)···H(1a)	0.97 (6)
P····C(14)	1.793 (6)	C(1)···H(1b)	1.04 (6)
C(2)···C(3)	1.390 (8)	C(1)···H(1c)	1.11 (9)
C(3)···C(4)	1.391 (8)	C(3)···H(3)	1.04 (5)
C(4)···C(5)	1.377 (10)	C(4)···H(4)	0.89 (7)
C(5)···C(6)	1.336 (10)	C(5)···H(5)	0.93 (5)
C(6)···C(7)	1.359 (8)	C(6)···H(6)	1.00 (5)
C(7)···C(2)	1.380 (7)	C(7)···H(7)	1.11 (5)
C(8)···C(9)	1.375 (10)	C(9)···H(9)	1.04 (6)
C(9)···C(10)	1.397 (12)	C(10)···H(10)	1.11 (9)
C(10)···C(11)	1.386 (10)	C(11)···H(11)	1.10 (7)
C(11)···C(12)	1.342 (12)	C(12)···H(12)	1.04 (7)
C(12)···C(13)	1.383 (12)	C(13)···H(13)	1.01 (7)
C(13)···C(8)	1.366 (8)	C(15)···H(15)	0.96 (6)
C(14)···C(15)	1.350 (7)	C(16)···H(16)	0.92 (6)
C(15)···C(16)	1.373 (9)	C(17)···H(17)	0.96 (7)
C(16)···C(17)	1.334 (13)	C(18)···H(18)	1.10 (8)
C(17)···C(18)	1.392 (9)	C(19)···H(19)	1.04 (5)
P····C'(8)	1.90 (6)	P····C'(14)	1.73 (5)
C'(8)···C'(9)	1.40 (7)	C'(14)···C'(15)	1.38 (8)
C'(9)···C'(10)	1.31 (6)	C'(15)···C'(16)	1.37 (9)
C'(10)···C'(11)	1.41 (6)	C'(16)···C'(17)	1.30 (11)
C'(11)···C'(12)	1.57 (9)	C'(17)···C'(18)	1.36 (8)
C'(12)···C'(13)	1.40 (8)	C'(18)···C'(19)	1.39 (6)
C'(13)···C'(8)	1.33 (6)	C'(19)···C'(14)	1.30 (7)
C(1)PC(2)	110.2 (3)°	H(1a)C(1)H(1b)	108 (4)°
C(1)PC(8)	109.2 (2)	H(1a)C(1)H(1c)	109 (5)
C(1)PC(14)	108.1 (2)	H(1b)C(1)H(1c)	108 (4)
C(2)PC(8)	109.8 (2)	C(2)C(3)H(3)	121 (3)
C(2)PC(14)	110.1 (2)	C(4)C(3)H(3)	121 (3)
C(8)PC(14)	109.6 (2)	C(3)C(4)H(4)	121 (4)
PC(2)C(3)	119.6 (4)	C(5)C(4)H(4)	118 (4)
PC(2)C(7)	120.4 (4)	C(4)C(5)H(5)	125 (3)
C(7)C(2)C(3)	119.9 (5)	C(6)C(5)H(5)	114 (3)
C(2)C(3)C(4)	117.8 (5)	C(5)C(6)H(6)	123 (3)
C(3)C(4)C(5)	120.1 (6)	C(7)C(6)H(6)	117 (3)
C(4)C(5)C(6)	121.4 (7)	C(6)C(7)H(7)	126 (2)
C(5)C(6)C(7)	119.7 (6)	C(2)C(7)H(7)	113 (2)
C(6)C(7)C(8)	120.9 (5)	C(8)C(9)H(9)	118 (3)
PC(8)C(9)	118.3 (4)	C(10)C(9)H(9)	123 (3)
PC(8)C(13)	122.2 (4)	C(9)C(10)H(10)	122 (3)
C(13)C(8)C(9)	119.5 (4)	C(11)C(10)H(10)	117 (3)
C(8)C(9)C(10)	118.9 (6)	C(10)C(11)H(11)	118 (2)
C(9)C(10)C(11)	120.5 (6)	C(12)C(11)H(11)	121 (2)
C(10)C(11)C(12)	119.6 (6)	C(11)C(12)H(12)	118 (4)
C(11)C(12)C(13)	120.2 (6)	C(13)C(12)H(12)	122 (4)
C(12)C(13)C(14)	121.1 (5)	C(12)C(13)H(13)	120 (3)
PC(14)C(15)	120.7 (4)	C(8)C(13)H(13)	119 (3)
PC(14)C(19)	121.7 (4)	C(14)C(15)H(15)	122 (3)
C(19)C(14)C(15)	117.5 (4)	C(16)C(15)H(15)	116 (3)
C(14)C(15)C(16)	122.0 (5)	C(15)C(16)H(16)	112 (4)
C(15)C(16)C(17)	121.8 (6)	C(17)C(16)H(16)	127 (4)
C(16)C(17)C(18)	118.6 (5)	C(16)C(17)H(17)	120 (3)
C(17)C(18)C(19)	120.6 (6)	C(18)C(17)H(17)	122 (3)
C(18)C(19)C(14)	119.4 (5)	C(17)C(18)H(18)	123 (4)
PC(1)H(1a)	118 (4)	C(19)C(18)H(18)	117 (4)
PC(1)H(1b)	106 (2)	C(18)C(19)H(19)	121 (3)
PC(1)H(1c)	108 (3)	C(14)C(19)H(19)	120 (3)
C(2)PC'(8)	109 (1)	C'(12)C'(13)C'(8)	121 (4)
C(2)PC'(14)	110 (1)	PC'(14)C'(15)	117 (3)
C'(8)PC'(14)	110 (2)	PC'(14)C'(19)	116 (3)
PC'(8)C'(9)	119 (3)	C'(19)C'(14)C'(15)	127 (4)
PC'(8)C'(13)	112 (3)	C'(14)C'(15)C'(16)	110 (3)
C'(13)C'(8)C'(9)	129 (3)	C'(15)C'(16)C'(17)	124 (5)
C'(8)C'(9)C'(10)	116 (4)	C'(16)C'(17)C'(18)	125 (4)
C'(9)C'(10)C'(11)	121 (4)	C'(17)C'(18)C'(19)	112 (3)
C'(10)C'(11)C'(12)	122 (4)	C'(18)C'(19)C'(14)	121 (3)
C'(11)C'(12)C'(13)	110 (4)		

37 and 68° respectively; in other words the 'handedness' of the propeller is inverted. In the other configuration (population 0.1), the handedness is the same as that at low temperature, phenyl rings II and III making angles of -20 and -49° respectively. As shown in Table 4, the shape and size of the cation are similar to those of the low-temperature phase. Relevant intermolecular contacts are listed in Table 5. The closest approach between tetrads packed side by side in the (100) layer occurs between hydrogen and nitrogen atoms. The N···H distances range from 2.48 to 2.80 Å. The closest N···N distance is 3.288 Å between N(51) and that related to this by a centre of symmetry at 0, $\frac{1}{2}$, $\frac{1}{2}$. The neighbours of $(CH_3PPh_3)^+$ will be discussed in the next section.

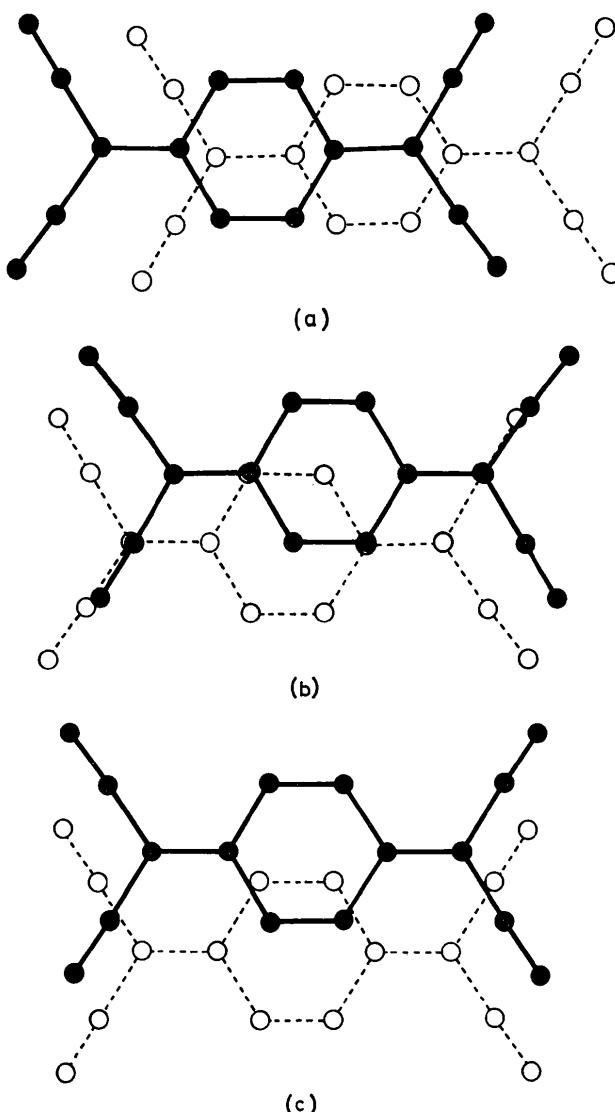


Fig. 4. Nearest neighbour overlap. (a) molecules I and I', II and I'; (b) molecules II' and II in the high-temperature modification; (c) molecules II' and II in the low-temperature modification.

Table 5. Relevant intermolecular distances

The C \cdots C and C \cdots N distances shorter than 3.5 Å and those involving the hydrogen atom shorter than 3.0 Å are listed. Key to symmetry operations

i	$-x, -y, -z$	vi	$1+x, y, 1+z$
ii	$-x, 1-y, -z$	vii	$1-x, 1-y, 1-z$
iii	$1+x, y, z$	viii	$x, y, 1+z$
iv	$1-x, -y, -z$	ix	$x, y, 1-z$
v	$1-x, 1-y, -z$		

Between TCNQ I and TCNQ II within a TCNQ tetrad

C(20) \cdots C(44)	3.358 (6) Å	C(20) \cdots C(46)	3.475 (6) Å
C(21) \cdots C(44)	3.455 (7)	C(21) \cdots C(46)	3.298 (8)
C(22) \cdots C(45)	3.380 (8)	C(23) \cdots C(36)	3.328 (8)
C(23) \cdots C(37)	3.474 (7)	C(24) \cdots C(36)	3.468 (8)
C(24) \cdots C(37)	3.338 (7)	C(25) \cdots C(36)	3.488 (7)
C(25) \cdots C(38)	3.422 (7)	C(28) \cdots C(47)	3.299 (7)
C(28) \cdots C(49)	3.376 (6)	C(29) \cdots C(47)	3.493 (6)
C(29) \cdots C(48)	3.439 (6)	C(30) \cdots C(49)	3.281 (6)

Between TCNQ I and TCNQ I' within a TCNQ tetrad

C(20) \cdots C(31 ⁱⁱ)	3.452 (8)	C(21) \cdots C(32 ⁱⁱ)	3.415 (7)
C(22) \cdots C(33 ⁱⁱ)	3.393 (7)	C(28) \cdots C(28 ⁱⁱ)	3.370 (9)
C(29) \cdots C(30 ⁱⁱ)	3.433 (8)		

Between TCNQ II and TCNQ II' of the two adjacent tetrads

C(39) \cdots C(48 ⁱⁱ)	3.479 (7)	C(40) \cdots N(50 ⁱⁱ)	3.418 (8)
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Between adjacent TCNQ columns

C(46) \cdots N(43 ^{vii})	3.428 (9)	H(46) \cdots N(43 ^{vii})	2.48 (8)
H(30) \cdots N(27 ^{vii})	2.69 (8)	H(22) \cdots N(35 ^{ix})	2.80 (6)
C(38) \cdots N(51 ^{ix})	3.500 (8)	H(38) \cdots N(51 ^{ix})	2.61 (7)
N(51) \cdots N(43 ^{vii})	3.358 (6)	N(51) \cdots N(51 ^{ix})	3.288 (11)
C(49) \cdots N(51 ^{ix})	3.366 (7)		

Between cation and TCNQ columns

Phenyl I			
C(7) \cdots N(50)	3.342 (7)	H(7) \cdots N(50)	2.49 (5)
H(7) \cdots N(34)	2.90 (6)	H(3) \cdots N(26 ⁱⁱ)	2.66 (5)
H(5) \cdots N(34 ^{ix})	2.88 (7)	H(6) \cdots N(35 ^{ix})	2.68 (6)

Phenyl II

H(9) \cdots N(34)	2.57 (6)	H(12) \cdots H(21 ⁱⁱ)	2.45 (10)
H(12) \cdots N(27 ⁱⁱ)	2.94 (8)	C(12) \cdots C(41 ^{vi})	3.499 (7)
C(12) \cdots N(43 ^{vi})	3.374 (8)	H(12) \cdots N(43 ^{vi})	2.99 (6)
C(13) \cdots N(43 ^{vi})	3.286 (7)	H(13) \cdots N(43 ^{vi})	2.81 (5)

Phenyl III

C(18) \cdots N(42 ⁱⁱⁱ)	3.387 (8)	*H(18) \cdots N(42 ⁱⁱⁱ)	2.91 (5)
C(19) \cdots N(42 ⁱⁱⁱ)	3.341 (9)	H(19) \cdots N(42 ⁱⁱⁱ)	2.83 (7)
*H(18) \cdots N(50 ^v)	2.91 (8)	H(15) \cdots C(40 ^{vi})	2.98 (6)
H(15) \cdots N(42 ^{vi})	2.85 (6)	H(16) \cdots C(48 ^{vi})	2.95 (6)
H(16) \cdots N(50 ^{vi})	2.90 (6)	*H(17) \cdots N(35 ^{vi})	2.92 (6)

Methyl

H(1b) \cdots N(34)	2.48 (6)	H(1c) \cdots N(50)	2.81 (7)
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Between cations

H(4) \cdots C(10 ^{iv})	2.89 (7)	C(3) \cdots H(10 ^{iv})	2.91 (9)
C(2) \cdots H(11 ^{iv})	2.91 (6)	C(3) \cdots H(11 ^{iv})	2.99 (7)
C(4) \cdots H(11 ^{iv})	2.95 (8)	C(5) \cdots H(11 ^{iv})	2.89 (8)
C(6) \cdots H(11 ^{iv})	2.79 (7)	C(7) \cdots H(11 ^{iv})	2.79 (6)

* The distance with an asterisk is shorter than the corresponding distance at low temperature.

Structural change with the transition

The volume of the unit cell contracts by about 2 Å³ above the transition temperature. This contraction is consistent with the observed phase diagram (Merkel *et al.*, 1965). Because of this very small change in cell dimensions, a single-crystal specimen transforms into the high-temperature phase without breaking. Thus the reported volume of the unit cell, 1751 Å³ seems to be somewhat incorrect (McPhail *et al.*, 1971). Changes in interplanar spacings between TCNQ moieties are shown in Table 6. The interplanar spacings within a tetrad increase by 0.08 and 0.04 Å, respectively, whereas that between two adjacent tetrads decreases by 0.02 Å. The overlapping mode in a tetrad is unaltered, but the nearest neighbour overlap between the tetrads changes as shown in Fig. 4.

Short intermolecular contacts between hydrogen atoms of the phenyl groups and nitrogen atoms of TCNQ moieties in the low-temperature modification are tabulated in Table 7. There are some short contacts less than the sum of van der Waals radii (2.70 Å) at room temperature. Such short contacts are similar to those observed in the crystals of methyltriphenoxyphosphonium bis-(1,2-dicyanoethylene-1,2-dithiolato)nickelate-(III) (Fritchie, 1966), tetraphenylarsonium 3-fluoro-1,1,4,5,5-pentacyano-2-azapentadienide (Palenik, 1966) and tetraphenylphosphonium (TCNQ)₂⁻ (Goldstein, Seff & Trueblood, 1968). As shown in Table 5, with one exception [H(9) \cdots N(34), 2.57 Å], no N \cdots H contact less than 2.70 Å is observed above the transition temperature. However, the rotation of the phenyl ring III increases the number of contacts less than 3.0 Å between hydrogen atoms of the cation with the nitrogen atoms of TCNQ II and II'. In Table 5 those distances with an asterisk are shorter than those in the low-temperature phase. In contrast to this, the rotation of the phenyl ring II simply gives rise to an increase in H \cdots N separations: they range from 2.41 to 2.80 Å in the low-temperature phase, whereas they are in the range 2.81~2.99 Å above the transition temperature apart from H(9) \cdots N(34) of 2.57 Å. These changes may be responsible for the elongation of the *a* axis by about 0.02 Å.

On the other hand the *b* and *c* axes contract by 0.18 and 0.11 Å respectively and the axial angle α decreases by about 3°. The plane of TCNQ moiety is inclined at an average angle of 64° with the *b* axis at room temperature; however, this angle of tilt increases to 67° above the transition temperature. Owing to this change in angle of tilt and the change in overlapping mode between adjacent tetrads, the increase in interplanar spacing within the tetrad may be alleviated, resulting in overall contraction of the volume.

Table 6. Interplanar spacings in both modifications

	Low-temperature phase	High-temperature phase	Difference
I \cdots I'	3.24 Å	3.32 Å	0.08 Å
I \cdots II'	3.22	3.26	0.04
II \cdots II'	3.57	3.55	-0.02

Table 7. Short intermolecular contacts between cation and TCNQ moieties in the low temperature phase (Tsuchiya, Marumo & Saito, 1973)

Phenyl II

C(13) ··· N(26 ^{III}) [*]	3.33 (1) Å	H(13) ··· N(26 ^{II})	2.65 (8) Å
H(12) ··· N(34 ^{IV})	2.41 (10)	C(9) ··· N(42 ^{VI})	3.38 (1)
H(9) ··· N(42 ^{VI})	2.80 (8)	C(10) ··· C(40 ^{VI})	3.48 (1)
H(10) ··· C(40 ^{VI})	2.96 (7)	C(10) ··· N(26 ^{VI})	3.62 (1)
H(10) ··· N(26 ^{VI})	2.51 (9)		

Phenyl III

H(19) ··· N(42 ^{II})	2.82 (7)	C(16) ··· H(37 ^{III})	2.86 (7)
C(15) ··· N(43 ^{VII})	3.30 (1)	H(15) ··· N(43 ^{VII})	2.51 (7)
H(16) ··· N(43 ^{VII})	2.86 (10)	H(16) ··· C(49 ^{VII})	2.83 (7)
H(16) ··· N(51 ^{VII})	2.64 (7)		

* Key to symmetry operation is given in Table 5.

The disorder in the conformation of $(CH_3PPh_3)^+$ ion contributes to the increase in the entropy of transition by $2.9 \text{ J}^\circ\text{K}^{-1}$. This is still 41% of the observed ΔS of transition (Kosaki *et al.*, 1970).

Values of anisotropic temperature factor indicate that the thermal vibration of the atoms is more vigorous above the transition temperature. Thus some other effect such as excitation of a new mode of thermal vibration might be responsible for the rest of ΔS . The increase in interplanar spacings in the tetrad may well account for the increase in magnetic susceptibility in the high-temperature phase. The decrease in interplanar spacing between the two tetrads is certainly related to the abrupt decrease in the line width of the exchange-narrowed line in the e.p.r. spectrum (Jones & Chesnut, 1963). The widening of interplanar spacing in the tetrads does not appear to favour the electrical conductivity. On the other hand, the decrease in interplanar spacing between the two adjacent tetrads, the approach of phenyl ring III to the nearest TCNQ II and II' and the decrease in the N(51) ··· N(51^{VII}) distance of 0.12 Å, might all be favourable to the electrical conductivity.

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